

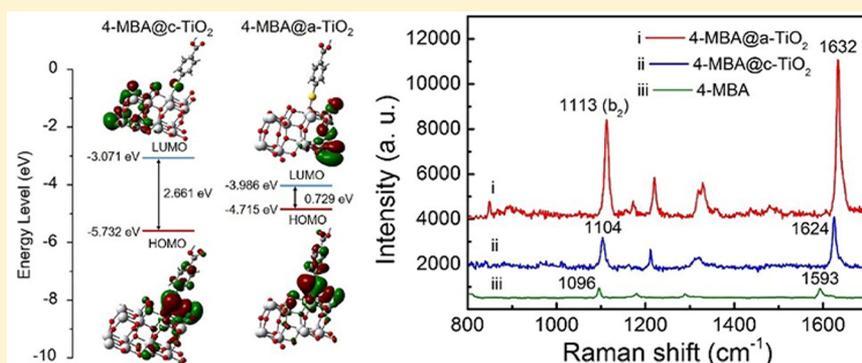
# Two-Dimensional Amorphous TiO<sub>2</sub> Nanosheets Enabling High-Efficiency Photoinduced Charge Transfer for Excellent SERS Activity

Xiaotian Wang,<sup>†,§</sup> Wenxiong Shi,<sup>‡,§</sup> Shaoxiong Wang,<sup>†,§</sup> Hewei Zhao,<sup>†</sup> Jie Lin,<sup>†</sup> Zhao Yang,<sup>†</sup> Mo Chen,<sup>†</sup> and Lin Guo<sup>\*,†,§</sup>

<sup>†</sup>School of Chemistry, Beijing Advanced Innovation Center for Biomedical Engineering, Key Laboratory of Bio-Inspired Smart Interfacial Science and Technology, Beihang University, 100191 Beijing, P. R. China

<sup>‡</sup>School of Material Science and Engineering, State Key Laboratory of Separation Membranes and Membrane Processes, Tianjin Polytechnic University, 300387 Tianjin, P. R. China

## Supporting Information



**ABSTRACT:** Substrate–molecule vibronic coupling enhancement, especially the efficient photoinduced charge transfer (PICT), is pivotal to the performance of nonmetal surface-enhanced Raman scattering (SERS) technology. Here, through developing novel two-dimensional (2D) amorphous TiO<sub>2</sub> nanosheets (a-TiO<sub>2</sub> NSs), we successfully obtained an ultrahigh enhancement factor of  $1.86 \times 10^6$ . Utilizing the Kelvin probe force microscopy (KPFM) technology, we found that these 2D a-TiO<sub>2</sub> NSs possessed more positive surface potential than their 2D crystalline counterpart (c-TiO<sub>2</sub> NSs). First-principles density functional theory (DFT) was used to further reveal that the low coordination number of surface Ti atoms and the large amount of surface oxygen defects endowed the 2D a-TiO<sub>2</sub> with high electrostatic potential, which allowed significant charge transfer from the adsorbed molecule to the 2D a-TiO<sub>2</sub> and facilitated the formation of a stable surface charge-transfer (CT) complex. Significantly, comparing with the 2D c-TiO<sub>2</sub>, the smaller band gap and higher electronic density of states (DOS) of the 2D a-TiO<sub>2</sub> effectively enhanced the vibronic coupling of resonances in the substrate–molecule system. The strong vibronic coupling within the CT complex obviously enhanced the PICT resonance and lead to the remarkable SERS activity of a-TiO<sub>2</sub> NSs. To the best of our knowledge, this is the first report on the remarkable SERS activity of 2D amorphous semiconductor nanomaterials, which may bring the cutting edge of development of stable and highly sensitive nonmetal SERS technology.

## INTRODUCTION

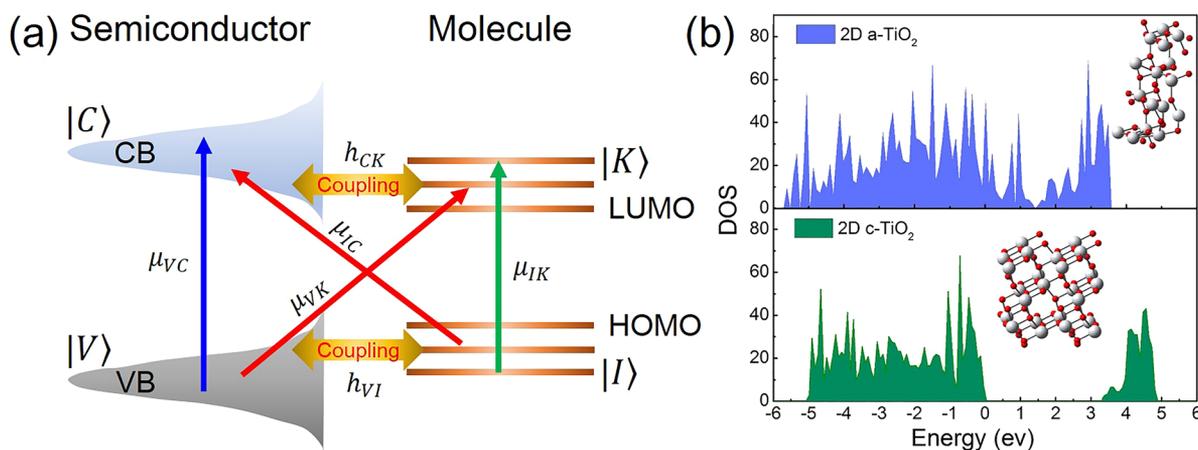
Surface-enhanced Raman scattering (SERS) spectroscopy owns the advantages of rapidity, high sensitivity, and nondestructiveness, which are extremely useful for label-free and fingerprint detection of multifold trace chemicals.<sup>1–4</sup> A few semiconductor nanomaterials have been proven to possess rather large enhancement factors exceeding  $10^4$  with high spectral stability, such as InAs/GaAs quantum dots,<sup>5</sup> CuTe nanocrystals,<sup>6</sup> Cu<sub>2</sub>O nanospheres,<sup>7</sup> and TiO<sub>2</sub> nanostructures,<sup>8,9</sup> in which photoinduced charge transfer (PICT) at the semiconductor–molecule interface is extremely important in Raman scattering enhancement.<sup>10,11</sup> Compared with metal nanostructures,<sup>12,13</sup> semiconductor nanostructures have excellent controllable properties,<sup>14–16</sup> such as exciton Bohr radius, band gap, n- and p-type doping, control of

stoichiometry, stability, stable exciton production, photoluminescence, resistance to degradation, and controllable geometry and crystalline phases through various synthetic techniques.<sup>17,18</sup> These results suggest that through engineering the crystal structure, energy band, and surface physicochemical properties to improve the PICT efficiency, semiconductor nanostructures may actually surpass metals in practical applications as SERS-active chemical sensors.

Two-dimensional (2D) semiconductor nanomaterials, which exhibit a flat surface with large specific surface area, were regarded as excellent candidates for fundamental studies of the SERS effect and its extensions to practical applications.<sup>19,20</sup>

Received: January 9, 2019

Published: March 21, 2019



**Figure 1.** (a) PICT transitions in the semiconductor–molecule system. (b) DOS calculation of 2D amorphous and crystalline TiO<sub>2</sub>.

Due to the extremely large specific surface area of 2D nanomaterials, the coordination number of the surface atoms is seriously insufficient, leading to easier combination of surface atoms with adsorbed molecules to form surface complexes,<sup>21</sup> which brings about stronger PICT resonance than the corresponding 3D structures.<sup>19,20</sup> However, because of the quantum effect of 2D materials,<sup>22</sup> the wide band gap and discontinuous band structure severely limit the vibronic coupling in the substrate–molecule system,<sup>23</sup> resulting in low PICT efficiency. In our previous work, we found that the numerous metastable electronic states of amorphous semiconductor nanostructures could serve as an intermediate level in the band gap, which benefited the interfacial charge transfer and resulted in the remarkable SERS activity of amorphous ZnO nanocages.<sup>24</sup> Therefore, combining the advantages of 2D nanomaterials with the amorphous structure, development of new 2D amorphous semiconductor nanomaterials might be a promising strategy to improve the interfacial PICT efficiency and boost the performance of semiconductor substrates accordingly.

Motivated by the above-mentioned exciting findings, we fabricated a new semiconductor SERS-active substrate, i.e., 2D amorphous TiO<sub>2</sub> nanosheets (a-TiO<sub>2</sub> NSs). This unique ultrathin nanostructure possesses unprecedented SERS activity with an enhancement factor (EF) up to  $1.86 \times 10^6$ , stronger than that of its 2D crystalline TiO<sub>2</sub> (c-TiO<sub>2</sub>) NS counterparts. Utilizing the Kelvin probe force microscopy (KPFM) technology, we found that these 2D a-TiO<sub>2</sub> NSs possessed more positive surface potential than their 2D crystalline counterpart (c-TiO<sub>2</sub> NSs). First-principles density functional theory (DFT) calculations further revealed that the lattice disorder and low coordination number of the surface Ti atoms would endow the 2D a-TiO<sub>2</sub> NSs with high surface potential and thus allowed a large amount of static charge transfer from the adsorbed molecule to the 2D a-TiO<sub>2</sub> and facilitated the formation of a stable surface charge-transfer (CT) complex. Significantly, comparing with the 2D c-TiO<sub>2</sub>, the smaller band gap and higher DOS of 2D a-TiO<sub>2</sub> can effectively improve the vibronic coupling of several resonances in the substrate–molecule system. The strong vibronic coupling within the CT complex will obviously enhance the PICT resonance and lead to the remarkable SERS activity of a-TiO<sub>2</sub> NSs. As far as we know, this is the first report of clear observation of the remarkable SERS activity from 2D amorphous semiconductors,

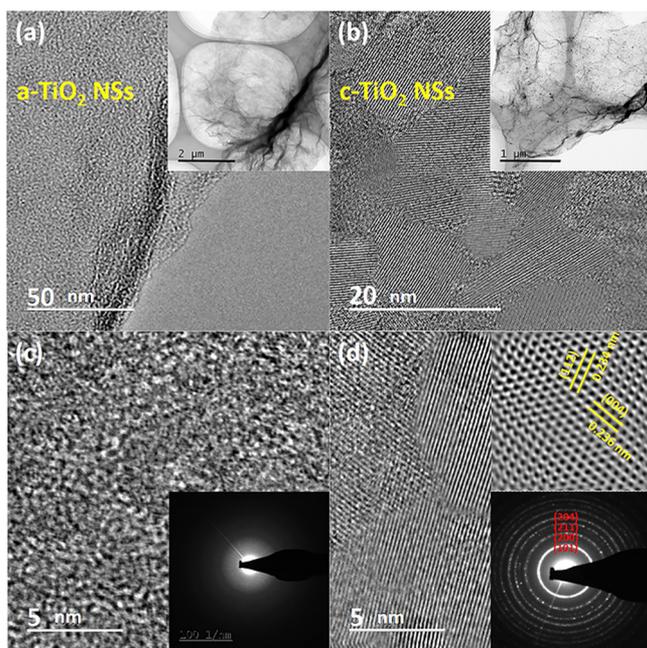
which may bring the cutting edge of development of stable and highly sensitive SERS technology.

## RESULTS AND DISCUSSION

According to the PICT mechanism, when the molecules are bound to the semiconductor surface via weak covalent bonds, the surface CT complex between the semiconductor and the molecule will be formed. This CT complex will achieve the strongest Raman enhancement once the excitation frequency resonates with the PICT transition.<sup>11,17,25</sup> Typically, the efficiency of PICT transition strongly depends on the intensity of vibronic coupling in the substrate–molecule system. When the energy levels of the molecules and the semiconductor substrate are coupled, PICT transitions from the semiconductor valence band (VB) states (*IV*) to the molecular excited states (*IK*) via the transition moment,  $\mu_{VK}$ , or from the molecular ground states (*II*) to the semiconductor conduction band (CB) states (*IC*) via the transition moment,  $\mu_{IC}$ , can borrow intensity either from exciton transitions  $\mu_{VC}$  or molecular transitions  $\mu_{IK}$ , through the Herzberg–Teller coupling term ( $h_{CK}$  or  $h_{VI}$ ), as shown in Figure 1a.

Therefore, developing a new strategy to effectively improve the vibronic coupling in the substrate–molecule system is the key issue to enhance the PICT resonance and the SERS effect. Here, we selected the TiO<sub>2</sub> as an ideal candidate for developing the high-performance SERS-active substrates due to its excellent biocompatibility, relatively high SERS activity, and good chemical stability among the reported semiconductors.<sup>26–28</sup> To investigate the vibronic coupling between the molecule and 2D amorphous semiconductor, we first calculated the electronic density of states (DOS) of 2D amorphous TiO<sub>2</sub> (a-TiO<sub>2</sub>). The DOS calculation of crystalline TiO<sub>2</sub> (c-TiO<sub>2</sub>) was also performed as a control. The DOS of 2D c- and a-TiO<sub>2</sub> were calculated at the density functional theory (DFT) plane-wave level based on the Vienna Ab-initio simulation package (VASP) combined with the projected-augmented wave method (PAW) (for details see Supporting Information S2). The results of DFT calculation showed that the 2D a-TiO<sub>2</sub> rendered smaller band gap and higher DOS than that of its crystal counterpart (Figure 1b). Although the DFT results of amorphous structures were qualitative assessments, it is sufficient to predict the energy band structure of 2D a-TiO<sub>2</sub> to enhance the vibronic coupling in the substrate–molecule system, suggesting a high-efficiency PICT process.

To achieve this, we designed and fabricated novel 2D amorphous TiO<sub>2</sub> nanosheets (a-TiO<sub>2</sub> NSs). The synthesis procedure follows our previous report with some modification.<sup>29</sup> Through manipulating the calcination temperature, the ultrathin a-TiO<sub>2</sub> NSs were successfully obtained, and the 2D crystalline TiO<sub>2</sub> nanosheets (c-TiO<sub>2</sub> NSs) were prepared as control samples (see **Materials and Methods** for details). Briefly, the highly uniform graphene was used as a template to synthesize the TiO<sub>2</sub> NSs@graphene precursor (Figure S1). After annealing at different temperatures in air, we successfully obtained the c- and a-TiO<sub>2</sub> NSs, respectively. The X-ray photoelectron spectroscopy (XPS) result also showed that there was no C element (C 1s: 285 eV) in 2D a-TiO<sub>2</sub> NSs (Figure S2), which suggested the graphene template had been removed via the heat treatment. It was observed from the XRD patterns (Figure S3) that the samples showed an amorphous phase when the calcination temperature was 400 °C. When the calcination temperature was elevated to 650 °C, the sharp diffraction lines with relatively high intensities indicated the high crystallinity of the anatase phase (JCPDS, no. 21-1272). Representative c- and a-TiO<sub>2</sub> NSs were further characterized by transmission electron microscopy (TEM) micrographs (Figure 2). It can be found that the as-prepared c- and a-TiO<sub>2</sub>



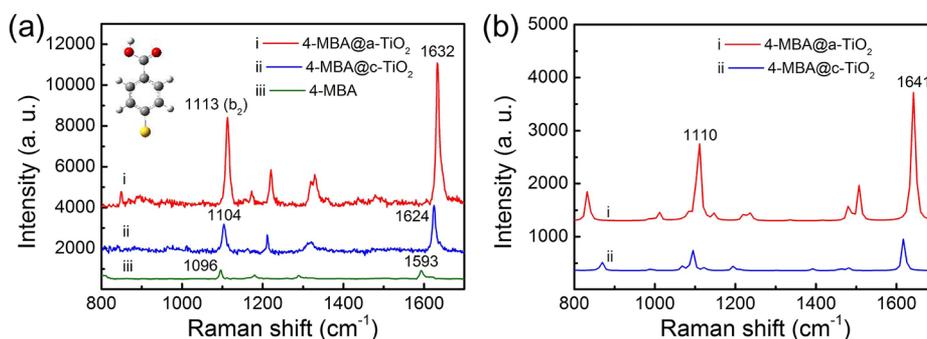
**Figure 2.** (a,b) TEM micrographs of a- and c-TiO<sub>2</sub> NSs. Inset: low-resolution images. (c,d) The HRTEM micrographs of a- and c-TiO<sub>2</sub> NSs. The insets in c,d are the corresponding SAED patterns of a- and c-TiO<sub>2</sub> NSs.

NSs showed flat ultrathin nanosheet morphology, indicating a successful inheritance of the graphene template shape. The low-resolution TEM micrographs in Figure 2a,b showed that the sizes of c- and a-TiO<sub>2</sub> NSs were around a few microns. The high-resolution TEM (HRTEM) micrograph (Figure 2c) showed that the amorphous a-TiO<sub>2</sub> NS structure was consistent with its corresponding selected area electron diffraction (SAED) pattern (the inset in Figure 2c). The inverse Fourier transform (FFT) image of c-TiO<sub>2</sub> NSs (inset in Figure 2d) displayed that the spacings of adjacent lattice fringe were 0.234 and 0.236 nm, respectively, which was consistent

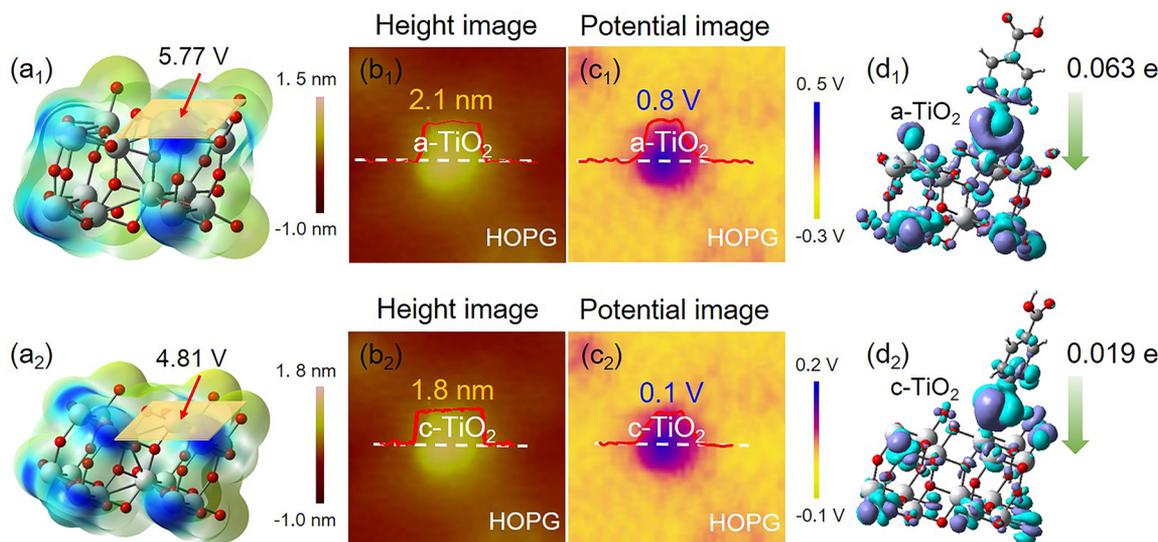
with the lattice constants of (112) and (004) planes of anatase TiO<sub>2</sub>. Moreover, the corresponding SAED patterns (the inset in Figure 2d) also clearly demonstrated the anatase phase of c-TiO<sub>2</sub> NSs, consistent with the XRD characterizations.

Moreover, Figure S4 showed the UV-vis-NIR spectra of the as-prepared c- and a-TiO<sub>2</sub> NSs. It can be seen that the absorption spectrum of the a-TiO<sub>2</sub> NSs possessed several broad absorption peaks and an obvious band tail, exhibiting an amorphous structure. A clear occurrence of absorption bands around 400 nm was ascribed to the interband excitation of the defect states,<sup>30</sup> owing to the well-known long-range disordered structure of amorphous materials. An abrupt change at 826 nm in the absorbance spectra suggested that the optical gap of the a-TiO<sub>2</sub> NSs was only 1.5 eV, much narrower than that of the c-TiO<sub>2</sub> NSs. The onset of optical absorption of the a-TiO<sub>2</sub> NSs was lowered to about 1.0 eV (~1230 nm). The broad absorption in the 850–1230 nm range can be attributed to the excitation of surface oxygen vacancy states.<sup>31</sup> The strong and broad absorption band within the whole visible region indicated a small band gap and a high DOS, which was consistent with the proposed DOS calculation (Figure 1b). The spectrum of the c-TiO<sub>2</sub> NSs showed an obviously decreased absorption above 400 nm, indicating a high crystallinity of the sample. The strong absorption below 390 nm was consistent with the band gap energy at 3.2 eV of anatase TiO<sub>2</sub>.

Based on the above results, it was reasonable that the relatively small band gap and high DOS allowed this novel 2D a-TiO<sub>2</sub> NS to effectively couple with the energy levels of the adsorbed molecule, resulting in the high-efficiency PICT resonance. Its flat surface and ultrathin structure also provided a promising approach for studying the SERS activity of 2D amorphous semiconductor nanomaterials. The standard probe molecule, i.e., 4-mercaptobenzoic acid (4-MBA), was used in our experiments. Figure 3a showed the control SERS experiments of c- and a-TiO<sub>2</sub> NSs modified by 4-MBA molecules (4-MBA@a-TiO<sub>2</sub> NSs and 4-MBA@c-TiO<sub>2</sub> NSs), respectively. It was demonstrated that a much stronger SERS effect occurred in a-TiO<sub>2</sub> NSs than its c-TiO<sub>2</sub> NS counterparts. First, compared with c-TiO<sub>2</sub> NSs, the features of the C=C stretching mode of the benzene ring (1593 cm<sup>-1</sup>) and ring-breathing mode coupled to the C-S stretch mode (1096 cm<sup>-1</sup>)<sup>31</sup> were remarkably enhanced in a-TiO<sub>2</sub> NSs. Second, compared with 4-MBA@c-TiO<sub>2</sub> NSs, more marked peak shifts of the probe molecules were found for 4-MBA@a-TiO<sub>2</sub> NSs, which was an indicator for strong semiconductor–molecule interaction.<sup>32</sup> As for 4-MBA, the intrinsic peak of 4-MBA (1096 and 1593 cm<sup>-1</sup>) shifted to 1113 and 1632 cm<sup>-1</sup> when adsorbed onto a-TiO<sub>2</sub> NSs and to 1104 and 1624 cm<sup>-1</sup> when adsorbed onto c-TiO<sub>2</sub> NSs, respectively. Third, we clearly observed the selective enhancement of nontotally symmetric lines (b<sub>2</sub> modes) in all c- and a-TiO<sub>2</sub> NSs, which was ascribed to the Herzberg–Teller contribution via a PICT mechanism.<sup>17</sup> For example, the C–H deformation mode at 1113 cm<sup>-1</sup> (assigned as a b<sub>2</sub> mode<sup>33</sup>) exhibited a large enhancement. Through collecting the SERS signals from ten different positions on 4-MBA molecule modified a-TiO<sub>2</sub> NSs (Figure S5), we found that these 2D amorphous substrates exhibited excellent spectra uniformity and repeatability. Moreover, according to the standard equation,<sup>34</sup> the EF was estimated to be approximately 1.86 × 10<sup>6</sup> for this a-TiO<sub>2</sub> NS (for the details, see Supporting Information S3). The remarkable SERS activity of a-TiO<sub>2</sub> NSs was further confirmed by using several



**Figure 3.** (a) Measured SERS spectra of 4-MBA ( $10^{-4}$  M) molecules adsorbed on c- and a-TiO<sub>2</sub> NSs and the normal Raman spectrum of pure 4-MBA solution (0.1 M) on the Si wafer. (b) Simulated SERS spectra of 4-MBA molecules adsorbed onto c- and a-TiO<sub>2</sub> NSs. Laser wavelength: 633 nm. Power: 0.8 mW. Lens: 50 $\times$  objective. Acquisition time: 10 s.



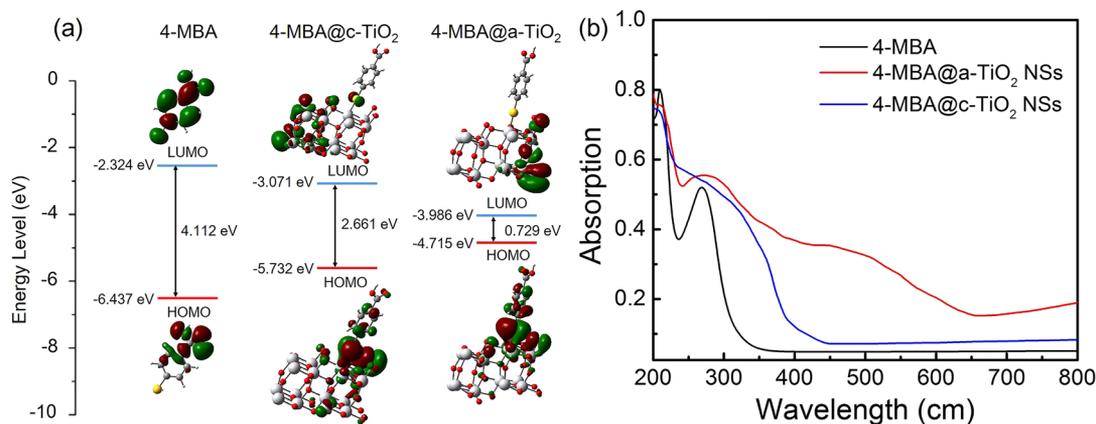
**Figure 4.** (a<sub>x</sub>) Calculated ESP at the cross-section surface of a- and c-TiO<sub>2</sub>. (b<sub>x</sub>) The AFM images of the single piece of a- and c-TiO<sub>2</sub> NSs on the HOPG substrate. (c<sub>x</sub>) The corresponding surface potential image of the single piece of a- and c-TiO<sub>2</sub> NSs on the HOPG substrate. (d<sub>x</sub>) The difference of charge distributions for 4-MBA@a- and c-TiO<sub>2</sub>, respectively. The cyan (purple) distribution illustrates electron depletion (accumulation). The electron transfer directions and values are also denoted. ( $x = 1$ ) a-TiO<sub>2</sub> and ( $x = 2$ ) c-TiO<sub>2</sub>.

standard probe molecules (Figure S7). It was clearly seen that the a-TiO<sub>2</sub> NSs had stronger Raman enhancement than their c-TiO<sub>2</sub> NS counterparts. Moreover, it was found that there was no big difference between the adsorption of 4-MBA molecules on c- and a-TiO<sub>2</sub> NSs through performing the UV-vis characterization of the 4-MBA ethanol solution after adsorption. Brunauer-Emmett-Teller (BET) N<sub>2</sub> adsorption-desorption curves for both c- and a-TiO<sub>2</sub> NSs were further characterized, respectively. The results showed that the specific surface areas (SSA) for c- and a-TiO<sub>2</sub> NSs were close to each other (3.37 m<sup>2</sup>/g for a-TiO<sub>2</sub> NSs and 3.53 m<sup>2</sup>/g for c-TiO<sub>2</sub> NSs, see Table S1 in Supporting Information), excluding the effects of differences in SSA between c- and a-TiO<sub>2</sub> NSs. Based on the above results, it was confirmed that the obvious SERS activities were derived from the unique surface physicochemical properties of a-TiO<sub>2</sub> NSs.

To understand the SERS effect of a-TiO<sub>2</sub> NSs, we carefully studied the surface CT complexes on 2D c- and a-TiO<sub>2</sub>, respectively. The electrostatic potential (ESP) and the difference of charge distributions in the two 4-MBA@c- and a-TiO<sub>2</sub> systems were investigated through DFT calculations. Briefly, the simulation 2D surface was modeled to be a slab consisting of two atomic layers (48 atoms). Typically, the 2D

a-TiO<sub>2</sub> model was obtained by molecular dynamics (MD) simulation to adjust the atomic coordination number, the bond length, and the atomic position in the 2D c-TiO<sub>2</sub> model (for details see Supporting Information S4). The mercapto molecules were chemically adsorbed onto both surfaces of 2D c- and a-TiO<sub>2</sub> models through the S-Ti bonds, which served as the interfacial charge transfer channels to facilitate the redistribution of the electron cloud around the TiO<sub>2</sub> and the molecule. Based on the above calculation models, the ESPs of 2D c- and a-TiO<sub>2</sub> were first calculated (for details see Supporting Information S5). The results in Figure 4a showed that the ESP at the cross-section surface of a-TiO<sub>2</sub> was 5.77 V, which was more positive than that of c-TiO<sub>2</sub> (4.81 V). More positive ESPs of 2D a-TiO<sub>2</sub> than its crystalline counterpart were ascribed to the low coordination number of surface Ti atoms and the large amount of surface oxygen defects, owing to the long-range disordered 2D amorphous structure.

To further confirm this, the surface potential of the c- and a-TiO<sub>2</sub> NSs was carefully characterized by Kelvin probe force microscopy (KPFM). KPFM technology was generally used to accurately characterize the surface potential and electron work function of materials, especially for 2D nanomaterials.<sup>35</sup> In our test, using highly oriented pyrolytic graphite (HOPG) as the



**Figure 5.** (a) Schematic diagram of the HOMO and LUMO of 4-MBA, 4-MBA@a-TiO<sub>2</sub>, and 4-MBA@c-TiO<sub>2</sub>. (b) Absorption spectra of 4-MBA, 4-MBA@a-TiO<sub>2</sub> NSs, and 4-MBA@c-TiO<sub>2</sub> NSs.

substrate, the surface potential of *c*- and *a*-TiO<sub>2</sub> NSs was accurately obtained by comparing the potential difference between the sample and the HOPG substrate. The AFM image and the corresponding surface potential image (tip lift: 10 nm) of a single piece of *c*- and *a*-TiO<sub>2</sub> nanosheets are shown in Figure 4b and Figure 4c, respectively. It can be clearly observed that the thickness of *c*- and *a*-TiO<sub>2</sub> NSs was around 2 nm, equivalent to a few monomolecular layers. The surface potentials of *c*- and *a*-TiO<sub>2</sub> NSs were 0.8 and 0.1 V, respectively, higher than that of the HOPG (the work function of HOPG: ~4.5 eV) substrate. Thus, the work functions of *c*- and *a*-TiO<sub>2</sub> NSs were 5.3 and 4.6 eV, respectively, which agreed well with the ESP calculations. Compared with the crystalline counterpart, 2D *a*-TiO<sub>2</sub> has more positive ESP and larger work function, which could strongly attract the electron cloud of the molecules and facilitated the interfacial charge transfer from the molecule to the *a*-TiO<sub>2</sub> NSs.

To investigate the interfacial charge transfer within the surface complex on both 2D *c*- and *a*-TiO<sub>2</sub> surfaces, the charge difference distributions for 4-MBA@*c*- and *a*-TiO<sub>2</sub> were calculated, respectively (Figure 4d). For details, please see Supporting Information S5. In general, the depleted region for electron density was around the C and S atoms in the benzene ring, and the accumulated region was near the Ti atoms in the TiO<sub>2</sub> substrate. The Hirshfeld population analysis<sup>36</sup> indicated that 0.063 e was transferred from 4-MBA to the 2D *a*-TiO<sub>2</sub>, while the amount of charge transferred from 4-MBA to the 2D *c*-TiO<sub>2</sub> cluster is 0.019 e. This result indicated that the 2D *a*-TiO<sub>2</sub> enabled more amount of charge to transfer from the 4-MBA molecules to 2D *a*-TiO<sub>2</sub>, suggesting strongly increased molecular polarizability.<sup>37,38</sup> To analyze the contribution of charge transfer to the molecular polarizability, the isotropic polarizabilities of probes 4-MBA@*c*- and *a*-TiO<sub>2</sub> were further calculated, respectively. The results revealed that the 2D *a*-TiO<sub>2</sub> could bring about much stronger amplification of the molecular polarization (see Supporting Information Table S3, 4-MBA@a-TiO<sub>2</sub>: 1449.66 Bohr<sup>3</sup>/atom, 4-MBA@c-TiO<sub>2</sub>: 1122.37 Bohr<sup>3</sup>/atom). Therefore, our calculations clearly demonstrated that the unique advantages of 2D *a*-TiO<sub>2</sub> effectively improved the interfacial charge transfer and thus magnified the molecular polarization. In addition, the calculated SERS spectra of the 4-MBA molecule chemically adsorbed onto the 2D *a*-TiO<sub>2</sub> surface exhibited larger peak shift and stronger intensity (Figure 3b), which were consistent with the experimental results. The changes in the intensity and

peak position between the simulation and measurement results can be attributed to the random molecule adsorption positions during the experiments and the differences among the ideal model settings in the simulations.

Besides improving the interfacial charge transfer and enhancing the molecular polarizability, we further found that the 2D *a*-TiO<sub>2</sub> greatly improved the energy level coupling between the adsorbed molecules and the 2D *a*-TiO<sub>2</sub> and effectively enhanced the PICT resonance within the surface CT complex. The DFT calculations showed that the HOMO and the LUMO for the 4-MBA molecules were distributed over the whole molecular structure (Figure 5a). When the 4-MBA molecules were adsorbed onto the 2D *c*- or *a*-TiO<sub>2</sub>, the surface CT complex was formed. The calculations showed that the LUMO of the 4-MBA@c-TiO<sub>2</sub> complex was still delocalized on the whole complex, while the LUMO of the 4-MBA@a-TiO<sub>2</sub> complex was mainly localized on the 2D *a*-TiO<sub>2</sub>. This clearly indicated that the 2D *a*-TiO<sub>2</sub> greatly modified the molecular orbitals of the adsorbed 4-MBA molecule after forming the surface CT complex. Significantly, it was found that the 2D *a*-TiO<sub>2</sub> could effectively reduce the band gap of the surface CT complex, which is favorable to the high-efficiency PICT resonance within the visible-light range. The calculation result clearly indicated that the band gap of the 4-MBA@a-TiO<sub>2</sub> complex was only 0.729 eV, smaller than that of the 4-MBA@c-TiO<sub>2</sub> complex (2.661 eV). As the HOMO–LUMO energy gap for 4-MBA was ~4.112 eV, the obviously reduced band gap of the 4-MBA@a-TiO<sub>2</sub> complex could be ascribed to the strong energy level coupling between the 4-MBA molecule and the 2D *a*-TiO<sub>2</sub>.

To investigate the intensity of the vibronic coupling within the surface CT complex, the band structures of *c*- and *a*-TiO<sub>2</sub> modified by 4-MBA were further calculated via the DFT calculation based on VASP. Figure S10 clearly showed that a comparison of the surface CT complex on *c*-TiO<sub>2</sub> (4-MBA@c-TiO<sub>2</sub>) with the 4-MBA@a-TiO<sub>2</sub> complex exhibited smaller band gap and more amount of resonant excited states within the whole visible light range, which allowed more possible thermodynamic PICT excitations at the low-energy level and hence effectively enhanced the Raman signals. This can be attributed to the strong energy level coupling between the 4-MBA molecule and the 2D *a*-TiO<sub>2</sub>, owing to the relatively small band gap and high DOS of 2D *a*-TiO<sub>2</sub>, which is consistent with our calculation in Figure 1b. In addition, the occurrence of strong PICT resonance in the 4-MBA@a-TiO<sub>2</sub>

complex was clearly confirmed through the UV–vis spectra measurements. Figure 5b showed the UV–vis spectra results of 4-MBA molecules, 4-MBA@a-TiO<sub>2</sub> NSs, and 4-MBA@c-TiO<sub>2</sub> NSs, respectively. It was observed that, compared with 4-MBA@c-TiO<sub>2</sub> NSs, 4-MBA@a-TiO<sub>2</sub> NSs exhibited an obvious red-shift and a broad absorption within the whole visible-light range, revealing the strong PICT resonance within the 4-MBA@a-TiO<sub>2</sub> NSs. In general, it was difficult to detect molecules with large energy gap because of the lack of the resonance Raman scattering (RRS) under visible light. Our findings suggest that these 2D a-TiO<sub>2</sub> NSs can provide a promising platform to generate efficient PICT resonance to realize the trace detection of molecules with large band gap.

Moreover, the wavelength-dependent SERS measurements were carried out (Figure S11). The SERS exhibited the strongest intensity under the 532 nm excitation compared with that under the 633 and 785 nm excitations. This can be ascribed to the strongest PICT resonance occurring under the 532 nm excitation, which agrees well with the relatively biggest absorption then (see UV–vis spectra in Figure 5b). Furthermore, the PICT pathways in both 4-MBA@c- and a-TiO<sub>2</sub> systems were carefully investigated (Figure S12). In brief, the HOMO energy levels of 4-MBA at  $-6.66$  eV were obtained by cyclic voltammetry (CV) (Figure S13). The HOMO–LUMO gap of 4-MBA was 4.43 eV from the UV–vis spectrum in Figure 5b. The band gap of a-TiO<sub>2</sub> NSs was 1.5 eV according to UV–vis–NIR spectra (Figure S4). The CB and VB of c-TiO<sub>2</sub> NSs were  $-4.2$  eV and  $-7.4$  eV below the vacuum level.<sup>39</sup> The valence bands of c- and a-TiO<sub>2</sub> NSs were measured via valence-band XPS<sup>40</sup> (Figure S14), which clearly indicated that the VB value toward the vacuum level for a-TiO<sub>2</sub> NSs was higher than that for c-TiO<sub>2</sub> NSs by 1.15 eV. As the VB of c-TiO<sub>2</sub> NSs was  $-7.4$  eV, the VB of a-TiO<sub>2</sub> NSs was  $-6.25$  eV. Therefore, the energy of charge-transfer transition  $\mu_{IC}$  (2.46 eV) is much higher than the excitation laser energy (633 nm, 1.96 eV) in the 4-MBA@c-TiO<sub>2</sub> system, impeding the thermodynamically allowed PICT process. However,  $\mu_{IC}$  (1.91 eV) is close to the excitation laser energy (633 nm, 1.96 eV) in the 4-MBA@a-TiO<sub>2</sub> system, which favors the efficient PICT between 4-MBA and a-TiO<sub>2</sub> NSs. Significantly, in the 4-MBA@a-TiO<sub>2</sub> system,  $\mu_{IC}$  can borrow the intensity from nearby exciton transition  $\mu_{VC}$  in a-TiO<sub>2</sub> NSs through the strong energy level coupling between the VB of a-TiO<sub>2</sub> NSs and the HOMO of 4-MBA. Thus, the a-TiO<sub>2</sub> NSs can exhibit much stronger SERS activity than c-TiO<sub>2</sub> NSs.

## CONCLUSION

In conclusion, through developing a novel ultrathin amorphous semiconductor nanomaterial, i.e., 2D a-TiO<sub>2</sub> NSs, an ultrahigh SERS activity (EF up to  $1.86 \times 10^6$ ) was successfully obtained in nonmetal nanomaterials. The underlying mechanism for this intriguing phenomenon was revealed via DFT simulations. Compared with the crystalline counterpart, more positive surface potential of 2D a-TiO<sub>2</sub> can strongly attract an electron cloud of molecules and endow a larger amount of interfacial charge transfer. Moreover, the relatively small band gap and high DOS enable the strong vibronic coupling to occur in the molecule@a-TiO<sub>2</sub> NS system. The above-mentioned mechanism can effectively enhance the PICT resonance and amplify the molecular Raman scattering. This progress will guide further development in design and fabrication of excellent SERS substrates and thrust the exploration of 2D amorphous

semiconductor nanomaterials for realizing ultrahigh-sensitivity SERS technology.

## ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/jacs.9b00029.

Figures, tables, synthesis methods and text giving characterization data for the 2D c- and a-TiO<sub>2</sub> NSs. DOS calculations of 2D c- and a-TiO<sub>2</sub> nanosheets, calculation of SERS enhancement factor (EF) of a-TiO<sub>2</sub> NSs, molecular dynamics (MD) simulation. XRD, absorption spectra and BET results for 2D c- and a-TiO<sub>2</sub> NSs, interface models and the contour of electrostatic potential (ESP) information (PDF)

## AUTHOR INFORMATION

### Corresponding Author

\*guolin@buaa.edu.cn

### ORCID

Wenxiang Shi: 0000-0002-7969-3780

Lin Guo: 0000-0002-6070-2384

### Author Contributions

<sup>§</sup>X. T. W., W. X. S., and S. X. W. contributed equally to this work.

### Notes

The authors declare no competing financial interest.

## ACKNOWLEDGMENTS

This project is supported by the National Natural Science Foundation of China (21875008, 51801007, 51532001).

## REFERENCES

- (1) Sanlessobrido, M.; Rodríguezlorenzo, L.; Lorenzoalbalde, S.; Gonzálezfernández, A.; Correaduarte, M. A.; Alvarezpuebla, R. A.; Lizmarzán, L. M. Label-free SERS detection of relevant bioanalytes on silver-coated carbon nanotubes: The case of cocaine. *Nanoscale* **2009**, *1*, 153–158.
- (2) Xu, L. J.; Zong, C.; Zheng, X. S.; Hu, P.; Feng, J. M.; Ren, B. Label-free detection of native proteins by surface-enhanced Raman spectroscopy using iodide-modified nanoparticles. *Anal. Chem.* **2014**, *86*, 2238–2245.
- (3) Driskell, J. D.; Tripp, R. A. Label-free SERS detection of microRNA based on affinity for an unmodified silver nanorod array substrate. *Chem. Commun.* **2010**, *46*, 3298–3300.
- (4) Schmidt, H. G.; Murphy, T.; Lucht, S.; Kronfeldt, H. D. Development of a SERS optode for the in-situ detection of chemicals in sea water. *Proceedings of SPIE-The International Society for Optical Engineering* **1999**, *3853*, 1015–1022.
- (5) Quagliano, L. G. Observation of Molecules Adsorbed on III-V Semiconductor Quantum Dots by Surface-Enhanced Raman Scattering. *J. Am. Chem. Soc.* **2004**, *126*, 7393–7398.
- (6) Li, W.; Zamani, R.; Gil, P. R.; Pelaz, B.; Ibáñez, M.; Cadavid, D.; Shavel, A.; Alvarezpuebla, R. A.; Parak, W. J.; Arbiol, J. CuTe Nanocrystals: Shape and Size Control, Plasmonic Properties, and Use as SERS Probes and Photothermal Agents. *J. Am. Chem. Soc.* **2013**, *135*, 7098–7101.
- (7) Jiang, L.; You, T.; Yin, P.; Shang, Y.; Zhang, D.; Guo, L.; Yang, S. Surface-enhanced Raman scattering spectra of adsorbates on Cu<sub>2</sub>O nanospheres: charge-transfer and electromagnetic enhancement. *Nanoscale* **2013**, *5*, 2784–2789.
- (8) Yang, L.; Gong, M.; Jiang, X.; Yin, D.; Qin, X.; Zhao, B.; Ruan, W. Investigation on SERS of different phase structure TiO<sub>2</sub> nanoparticles. *J. Raman Spectrosc.* **2015**, *46*, 287–292.

- (9) Wang, Y.; Yan, C.; Chen, L.; Zhang, Y.; Yang, J. Controllable Charge Transfer in Ag-TiO<sub>2</sub>(2) Composite Structure for SERS Application. *Nanomaterials* **2017**, *7*, 159.
- (10) Persson, B. N.; Zhao, K.; Zhang, Z. Chemical contribution to surface-enhanced Raman scattering. *Phys. Rev. Lett.* **2006**, *96*, 207401.
- (11) Wang, X.; Shi, W.; She, G.; Mu, L. Using Si and Ge nanostructures as substrates for surface-enhanced Raman scattering based on photoinduced charge transfer mechanism. *J. Am. Chem. Soc.* **2011**, *133*, 16518–16523.
- (12) Wang, X.; Ma, G.; Li, A.; Yu, J.; Yang, Z.; Lin, J.; Li, A.; Han, X.; Guo, L. Composition-adjustable Ag-Au substitutional alloy microcages enabling tunable plasmon resonance for ultrasensitive SERS. *Chem. Sci.* **2018**, *9*, 4009–4015.
- (13) Tao, L.; Chen, K.; Chen, Z.; Cong, C.; Qiu, C.; Chen, J.; Wang, X.; Chen, H.; Yu, T.; Xie, W.; Deng, S.; Xu, J. B. 1T' Transition Metal Telluride Atomic Layers for Plasmon-Free SERS at Femtomolar Levels. *J. Am. Chem. Soc.* **2018**, *140*, 8696–8704.
- (14) Wang, X.; Shi, W.; She, G.; Mu, L. Surface-Enhanced Raman Scattering (SERS) on transition metal and semiconductor nanostructures. *Phys. Chem. Chem. Phys.* **2012**, *14*, 5891–5901.
- (15) Jin, L.; She, G.; Wang, X.; Mu, L.; Shi, W. Enhancing the SERS performance of semiconductor nanostructures through a facile surface engineering strategy. *Appl. Surf. Sci.* **2014**, *320*, 591–595.
- (16) Zheng, Z.; Cong, S.; Gong, W.; Xuan, J.; Li, G.; Lu, W.; Geng, F.; Zhao, Z. Semiconductor SERS enhancement enabled by oxygen incorporation. *Nat. Commun.* **2017**, *8*, 1993.
- (17) Lombardi, J. R.; Birke, R. L. Theory of Surface-Enhanced Raman Scattering in Semiconductors. *J. Phys. Chem. C* **2014**, *118*, 11120–11130.
- (18) Kaur, K.; Singh, C. V. Amorphous TiO<sub>2</sub> as a photocatalyst for hydrogen production: a DFT study of structural and electronic properties. *Energy Procedia* **2012**, *29*, 291–299.
- (19) Ling, X.; Fang, W.; Lee, Y. H.; Araujo, P. T.; Zhang, X.; Rodriguez-Nieva, J. F.; Lin, Y.; Zhang, J.; Kong, J.; Dresselhaus, M. S. Raman enhancement effect on two-dimensional layered materials: graphene, h-BN and MoS<sub>2</sub>. *Nano Lett.* **2014**, *14*, 3033–3040.
- (20) Muehlethaler, C.; Considine, C. R.; Menon, V.; Lin, W.-C.; Lee, Y.-H.; Lombardi, J. R. Ultrahigh Raman Enhancement on Monolayer MoS<sub>2</sub>. *ACS Photonics* **2016**, *3*, 1164–1169.
- (21) Cong, S.; Yuan, Y.; Chen, Z.; Hou, J.; Yang, M.; Su, Y.; Zhang, Y.; Li, L.; Li, Q.; Geng, F.; Zhao, Z. Noble metal-comparable SERS enhancement from semiconducting metal oxides by making oxygen vacancies. *Nat. Commun.* **2015**, *6*, 7800.
- (22) Cudazzo, P.; Sponza, L.; Giorgetti, C.; Reining, L.; Sottile, F.; Gatti, M. Exciton Band Structure in Two-Dimensional Materials. *Phys. Rev. Lett.* **2016**, *116*, 066803.
- (23) Lin, J. H.; Zhang, H.; Cheng, X. L.; Miyamoto, Y. Two-dimensional wide-band-gap nitride semiconductors: Single-layer 1T-XN<sub>2</sub> (X = S, Se, and Te). *Phys. Rev. B: Condens. Matter Mater. Phys.* **2016**, *94*, 195404.
- (24) Wang, X.; Shi, W.; Jin, Z.; Huang, W.; Lin, J.; Ma, G.; Li, S.; Guo, L. Remarkable SERS Activity Observed from Amorphous ZnO Nanocages. *Angew. Chem., Int. Ed.* **2017**, *56*, 9851–9855.
- (25) Alessandri, I.; Lombardi, J. R. Enhanced Raman Scattering with Dielectrics. *Chem. Rev.* **2016**, *116*, 14921–14981.
- (26) Scanlon, D. O.; Dunnill, C. W.; Buckeridge, J.; Shevlin, S. A.; Logsdail, A. J.; Woodley, S. M.; Catlow, C. R. A.; Powell, M. J.; Palgrave, R. G.; Parkin, I. P.; Watson, G. W.; Keal, T. W.; Sherwood, P.; Walsh, A.; Sokol, A. A. Band alignment of rutile and anatase TiO<sub>2</sub>. *Nat. Mater.* **2013**, *12*, 798–801.
- (27) Musumeci, A.; Gosztola, D.; Schiller, T.; Dimitrijevic, N. M.; Mujica, V.; Martin, D.; Rajh, T. SERS of Semiconducting Nanoparticles (TiO<sub>2</sub> Hybrid Composites). *J. Am. Chem. Soc.* **2009**, *131*, 6040–6041.
- (28) Bontempi, N.; Vassalini, I.; Alessandri, I. All-dielectric core/shell resonators: From plasmon-free SERS to multimodal analysis. *J. Raman Spectrosc.* **2018**, *49*, 943–953.
- (29) Zhao, H.; Zhu, Y.; Li, F.; Hao, R.; Wang, S.; Guo, L. A Generalized Strategy for Synthesis of Large-size Ultrathin Two-dimensional Metal Oxide Nanosheets. *Angew. Chem., Int. Ed.* **2017**, *56*, 8766–8770.
- (30) Iijima, K.; Goto, M.; Enomoto, S.; Kunugita, H.; Ema, K.; Tsukamoto, M.; Ichikawa, N.; Sakama, H. Influence of oxygen vacancies on optical properties of anatase TiO<sub>2</sub> thin films. *J. Lumin.* **2008**, *128*, 911–913.
- (31) Wang, Y.; Sun, Z.; Hu, H.; Jing, S.; Zhao, B.; Xu, W.; Zhao, C.; Lombardi, J. R. Raman scattering study of molecules adsorbed on ZnS nanocrystals. *J. Raman Spectrosc.* **2007**, *38*, 34–38.
- (32) Wang, Z.; Rothberg, L. J. Origins of blinking in single-molecule Raman spectroscopy. *J. Phys. Chem. B* **2005**, *109*, 3387–3391.
- (33) Jiang, X.; Li, X.; Jia, X.; Li, G.; Wang, X.; Wang, G.; Li, Z.; Yang, L.; Zhao, B. Surface-Enhanced Raman Scattering from Synergistic Contribution of Metal and Semiconductor in TiO<sub>2</sub>/MBA/Ag(Au) and Ag(Au)/MBA/TiO<sub>2</sub> Assemblies. *J. Phys. Chem. C* **2012**, *116*, 14650–14655.
- (34) Orendorff, C. J.; Gole, A.; Sau, T. K.; Murphy, C. Murphy, Surface-Enhanced Raman Spectroscopy of Self-Assembled Monolayers: Sandwich Architecture and Nanoparticle Shape Dependence. *Anal. Chem.* **2005**, *77*, 3261–3266.
- (35) Lee, S. H.; Lee, S. W.; Oh, T.; Petrosko, S. H.; Mirkin, C. A.; Jang, J. W. Direct observation of plasmon-induced interfacial charge separation in metal/semiconductor hybrid nanostructures by measuring surface potentials. *Nano Lett.* **2018**, *18*, 109.
- (36) Hirshfeld, F. L. Bonded-atom fragments for describing molecular charge densities. *Theoretica Chimica Acta* **1977**, *44*, 129–138.
- (37) Zebarjadi, M.; Joshi, G.; Zhu, G.; Yu, B.; Minnich, A.; Lan, Y.; Wang, X.; Dresselhaus, M.; Ren, Z.; Chen, G. Power Factor Enhancement by Modulation Doping in Bulk Nanocomposites. *Nano Lett.* **2011**, *11*, 2225.
- (38) Alessandri, I. Enhancing Raman scattering without plasmons: unprecedented sensitivity achieved by TiO<sub>2</sub> shell-based resonators. *J. Am. Chem. Soc.* **2013**, *135*, 5541.
- (39) Hou, W.; Cronin, S. B. A Review of Surface Plasmon Resonance-Enhanced Photocatalysis. *Adv. Funct. Mater.* **2013**, *23*, 1612–1619.
- (40) Chen, X.; Liu, L.; Yu, P. Y.; Mao, S. S. Increasing Solar Absorption for Photocatalysis with Black Hydrogenated Titanium Dioxide Nanocrystals. *Science* **2011**, *331*, 746–750.